# Geometry and Computability of zero-temperature measures

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# Outline of the talk

- Pressure, equilibrium states and zero-temperature measures
- 2 Existence of zero-temperature measures
- Computability of zero-temperature measures and residual entropy

Setup: Let (X, d) be a compact metric space and let  $f : X \to X$  be continuous. We endow

 $\mathcal{M} = \{\mu : f - \text{invariant Borel probability measure on } X\}$ 

with the weak\* topology.

 $\implies \mathcal{M}$  a compact, convex, metrizable.

Let  $\mathcal{M}_E \subset \mathcal{M}$  be the subset of ergodic measures, that is, if  $f^{-1}(A) = A$  then  $\mu(A) \in \{0, 1\}$ .

We denote by  $h_{\mu}(f)$  the measure-theoretic entropy of  $\mu$ .

Let  $C(X, \mathbb{R})$  denote the space of continuous functions  $\phi : X \to \mathbb{R}$  endowed with the supremums norm.

#### Definition

The topological pressure  $P_{top}$  :  $C(X, \mathbb{R}) \to \mathbb{R}$  is defined by

$$P_{ ext{top}}(\phi) = \sup_{\mu \in \mathcal{M}} h_{\mu}(f) + \int \phi \, d\mu$$

A measure  $\mu$  is an *equilibrium state* of  $\phi$  if  $P_{top}(\phi) = h_{\mu}(f) + \int \phi \, d\mu$ . The topological pressure of the constant zero potential is called the topological entropy  $h_{top}(f)$  of f.

We assume that the entropy map  $\mu \mapsto h_{\mu}(f)$  is upper semi-continuous which guarantees that for each potential  $\phi$  there exists at least one (ergodic) equilibrium state.

Let  $\phi \in C(X, \mathbb{R})$  such that for all  $\beta \geq 0$  the potential  $\beta \phi$  has a unique equilibrium state  $\mu_{\beta} = \mu_{\beta\phi}$ . We think of  $\beta = \frac{1}{T}$  as the inverse temperature (*T* being the temperature) of the system.

#### Definition

If the limit  $\mu_{\infty,\phi} = \lim_{\beta \to \infty} \mu_{\beta}$  (in the weak\* topology) exists we call  $\mu_{\infty,\phi}$  the zero-temperature measure of  $\phi$ . If  $\mu_{\infty,\phi}$  does not exist we consider the set of of accumulation points of  $\mu_{\beta}$  as  $\beta \to \infty$  which we call the ground state set  $GS(\phi)$  of  $\phi$ .

#### Theorem (Folklore)

Suppose  $\mu \in GS(\phi)$ . Then

(i) 
$$\int \phi \, d\mu = \sup \{ \int \phi \, d\nu : \nu \in \mathcal{M} \} \stackrel{\text{def}}{=} b_{\phi}.$$

(ii) 
$$h_{\mu}(f) = \sup\{h_{\nu}(f) : \int \phi \, d\nu = b_{\phi}\} \stackrel{\text{def}}{=} h_{\infty,\phi}.$$

The quantity  $h_{\infty,\phi}$  is called the residual entropy of  $\phi$ , i.e., the maximal entropy at zero temperature.

Let  $d \ge 2$ ,  $\mathcal{A} = \{0, \ldots, d-1\}$  be a finite alphabet and let A be a transition matrix. We define  $X = X_A = \{(x_i)_{i \in \mathbb{N}_0} : x_i \in \mathcal{A}, A_{x_i, x_{i+1}} = 1\}$ . Let  $f : X \to X$  be the shift map. We say f is a subshift of finite type. In the following we assume that f is transitive.

Denote by  $LC(X, \mathbb{R}) = \bigcup_k LC_k(X, \mathbb{R})$  the locally constant potentials, where  $LC_k(X, \mathbb{R})$  are the potentials constant on cylinders of length k.

#### Theorem (Brémont, Nonlinearity, 2003)

If  $\phi \in LC(X, \mathbb{R})$  then the zero-temperature measure  $\mu_{\infty,\phi}$  exists.

#### Theorem (W., Yang, TAMS, 2019)

Let  $k \in \mathbb{N}$ . Then there exist a partition of  $LC_k(X, \mathbb{R})$  into finitely many convex cones  $U_1, \ldots, U_{\ell_1}, U_{\ell_1+1}, \ldots, U_N$  and k-elementary points  $x_1, \ldots, x_\ell$  such that:

- (i) U<sub>1</sub>,...,U<sub>ℓ1</sub> are open, and U<sub>1</sub> ∪ ... ∪ U<sub>ℓ</sub> = LC<sub>k</sub>(X, ℝ). Further, the orbit of x<sub>i</sub> is the unique maximizing periodic orbit of φ. For each i = 1,..., ℓ, for all φ ∈ U<sub>i</sub>, we have μ<sub>∞,φ</sub> = μ<sub>xi</sub>, i.e., the unique invariant measure supported on the periodic orbit x<sub>i</sub>.
- (ii) If  $\phi \in U_{\ell+1} \cup \ldots \cup U_N$  then  $\mu_{\infty,\phi}$  is either non-ergodic or has positive entropy or both.

## Basics from Computable Analysis:

#### Definition

Let  $m \in \mathbb{N}$  and  $x \in \mathbb{R}^m$ . An *oracle* of x is a function  $\psi : \mathbb{N} \to \mathbb{Q}^m$  such that  $\|\psi(n) - x\| < 2^{-n}$ . Moreover, x is computable if there is a Turing Machine (a computer program for our purposes)  $\psi$  which is an oracle of x.

Basic Facts:

- (i) Rational numbers, algebraic numbers, and some transcendental numbers such as e and  $\pi$  are computable real numbers.
- (ii) There are only countably many computable points in  $\mathbb{R}^m$ .

#### Definition

Let  $S \subset \mathbb{R}^m$ . A function  $g : S \to \mathbb{R}$  is *computable* if there is a Turing machine  $\chi$  so that for any  $x \in S$  and any oracle  $\psi$  for x,  $\chi(\psi, n)$  is a rational number so that  $|\chi(\psi, n) - g(x)| < 2^{-n}$ .

#### Definition

Let  $S \subset \mathbb{R}^m$ . A function  $g: S \to \mathbb{R}$  is *upper semi-computable* if there is a Turing machine  $\chi$  so that for any  $n \in \mathbb{N}$  and  $x \in S$ , and any oracle  $\psi$  for x,  $\chi(\psi, n) = q_n \in \mathbb{Q}$  such that  $q_n \downarrow g(x)$  as  $n \to \infty$ . Analogously one defines lower semi-computable functions.

#### Lemma

A function  $g: S \to \mathbb{R}$  is computable if and only if g is lower and upper semi-computable.

Similar notions of computability exist for many mathematical objects (subsets of Euclidian space, shift spaces, measure spaces, etc.) and functions.

## A Computability result for zero-temperature measures:

#### Theorem (Burr, W., Nonlinearity 2021)

Let  $k \in \mathbb{N}$ . Let  $\mathcal{O}_k$  be the set of potentials  $\phi \in LC_k(X, \mathbb{R})$  such that the zero-temperature measure  $\mu_{\infty,\phi}$  is supported on a periodic orbit. Then

- (i) Then  $\mathcal{O}_k \ni \phi \mapsto \mu_{\infty,\phi}$  is computable.
- (ii) The set  $\mathcal{O}_k$  is a lower computable open and dense subset of  $LC_k(X, \mathbb{R})$ .
- (iii) If  $\phi_0 \in LC_k(X, \mathbb{R}) \setminus \mathcal{O}_k$ , then  $\mu_{\infty,\phi_0}$  is not computable.
- (iv) If the cylinder length k is not given in advance, then  $\mu_{\infty,\phi_0}$  is not computable for any  $\phi_0 \in LC(X, \mathbb{R})$ .

# A Computability result for the residual entropy:

#### Theorem (Burr, W., Nonlinearity 2021)

Let  $f : X \to X$  be a transitive subshift of finite type. Then the function  $\phi \mapsto h_{\infty,\phi}$  is upper semi-computable, but not computable on  $C(X,\mathbb{R})$ . Moreover, the map  $\phi \mapsto h_{\infty,\phi}$  is continuous at  $\phi_0$  if and only if  $h_{\infty,\phi_0} = 0$ .